ONE-LOOP NEUTRINO MASS IN $SU(5)^*$

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*I. Doršner, S. Fajfer, N. Košnik, work in progress.

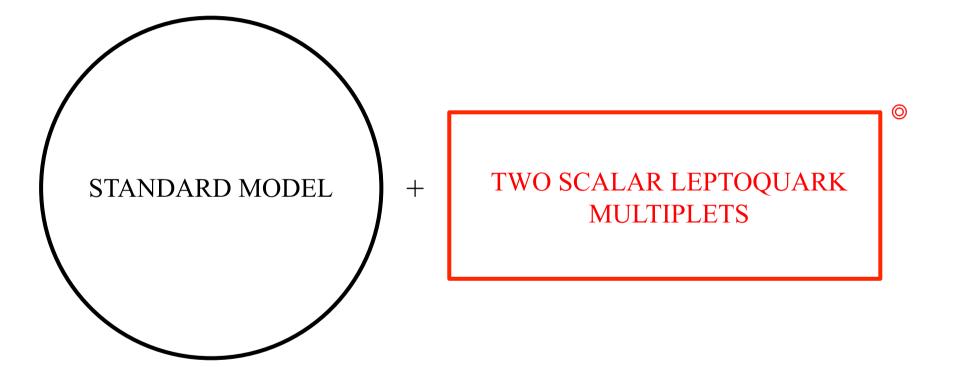
OUTLINE

•ONE-LOOP NEUTRINO MASS WITH LEPTOQUARKS

•ONE-LOOP NEUTRINO MASS MECHANISM IN SU(5)

CONCLUSIONS

ONE-LOOP NEUTRINO MASS MECHANISM



[©]C.-K. Chua, X.-G. He, W.-Y. P. Hwang, Phys. Lett. B479 (2000).

SCALAR LEPTOQUARKS[©]

LEPTOQUARK (LQ) MULTIPLETS:

 $(\overline{\mathbf{3}}, \mathbf{3}, 1/3)$ $(\mathbf{3}, \mathbf{2}, 7/6)$ $(\overline{\mathbf{3}}, \mathbf{2}, 1/6)$ $(\overline{\mathbf{3}}, \mathbf{1}, 1/3)$ $(\overline{\mathbf{3}}, \mathbf{1}, 4/3)$

$$({\bf \overline{3}},{\bf 1},-2/3)$$

[©]I. Doršner, S. Fajfer, A. Greljo, J.F. Kamenik, N. Košnik, Phys. Rept. 641 (2016).

SCALAR LEPTOQUARKS

LQ NOMENCLATURE[©]:

$$S_3 \equiv (\overline{\mathbf{3}}, \mathbf{3}, 1/3)$$

$$R_2 \equiv (\mathbf{3}, \mathbf{2}, 7/6)$$

$$\tilde{R}_2 \equiv ({\bf 3}, {\bf 2}, 1/6)$$

$$S_1 \equiv (\overline{\mathbf{3}}, \mathbf{1}, 1/3)$$

$$\tilde{S}_1 \equiv (\overline{\mathbf{3}}, \mathbf{1}, 4/3)$$

$$S_1 \equiv (\overline{\mathbf{3}}, \mathbf{1}, -2/3)$$

[©]W. Büchmuller *et al.*, Phys. Lett. B 191, 442 (1987).

SCALAR LEPTOQUARKS VS. *v* **MASS**

v MASS LQs: \tilde{R}_2 + (S_3 \lor S_1)

$$S_3 \equiv (\overline{\mathbf{3}}, \mathbf{3}, 1/3)$$

$$R_2 \equiv (\mathbf{3}, \mathbf{2}, 7/6)$$

$$\tilde{R}_2 \equiv ({\bf 3}, {\bf 2}, 1/6)$$

$$S_1 \equiv (\overline{\mathbf{3}}, \mathbf{1}, 1/3)$$

$$\tilde{S}_1 \equiv (\overline{\mathbf{3}}, \mathbf{1}, 4/3)$$

$$\bar{S}_1 \equiv (\overline{\mathbf{3}}, \mathbf{1}, -2/3)$$

SCALAR LEPTOQUARKS VS. *p* **DECAY**

p DECAY LQs: ($S_3,\, ilde R_2,\,S_1,\, ilde S_1,\,ar S_1$)

 $S_3 \equiv (\overline{\mathbf{3}}, \mathbf{3}, 1/3)$

$$R_2 \equiv (\mathbf{3}, \mathbf{2}, 7/6)$$

$$ilde{R}_2\equiv~~({f 3},{f 2},1/6)$$

$$S_1 \equiv (\overline{\mathbf{3}}, \mathbf{1}, 1/3)$$

$$ilde{S}_1 \equiv (\overline{\mathbf{3}}, \mathbf{1}, 4/3)$$

$$ar{S}_1 \equiv \ (\overline{\mathbf{3}}, \mathbf{1}, -2/3)$$

v MASS VS. p DECAY

$$S_3 \equiv (\overline{\mathbf{3}}, \mathbf{3}, 1/3)$$

$$R_2 \equiv (\mathbf{3}, \mathbf{2}, 7/6)$$

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$$\tilde{S}_1 \equiv (\overline{\mathbf{3}}, \mathbf{1}, 4/3)$$

$$\bar{S}_1 \equiv (\overline{\mathbf{3}}, \mathbf{1}, -2/3)$$

A(NOTHER) WORD ABOUT NOMENCLATURE

$$\mathcal{L} \supset -\tilde{y}_{2\,ij}^{RL} \bar{d}_R^i \tilde{R}_2^a \epsilon^{ab} L_L^{j,b}$$

a, b (= 1, 2) are SU(2) indices i, j (= 1, 2, 3) are flavor indices $\tilde{y}_2^{RL} \equiv$ Yukawa coupling matrix

A(NOTHER) WORD ABOUT NOMENCLATURE

$$\mathcal{L} \supset -\tilde{y}^{RL}_{2\,ij} \bar{d}^i_R \tilde{R}^a_2 \epsilon^{ab} L^{j,b}_L$$



$$\mathcal{L} \supset -\tilde{y}_{2\,ij}^{RL} \bar{d}_R^i e_L^j \tilde{R}_2^{2/3} + (\tilde{y}_2^{RL} U)_{ij} \bar{d}_R^i \nu_L^j \tilde{R}_2^{-1/3}$$

a, *b* (= 1, 2) are *SU*(2) indices

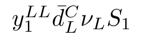
i, j (= 1, 2, 3) are flavor indices

 $\tilde{y}_2^{RL} \equiv$ Yukawa coupling matrix

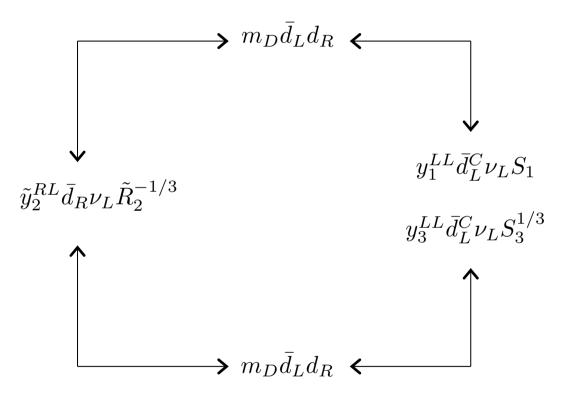
 $U \equiv$ Pontecorvo-Maki-Nakagawa-Sakata unitary mixing matrix

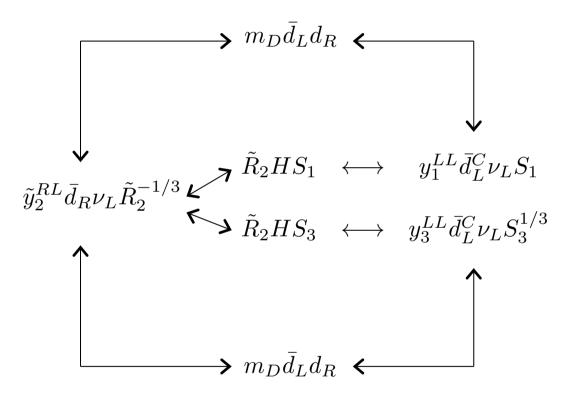
 $\tilde{y}_2^{RL} \bar{d}_R \nu_L \tilde{R}_2^{-1/3}$

 $\tilde{y}_2^{RL} \bar{d}_R \nu_L \tilde{R}_2^{-1/3}$

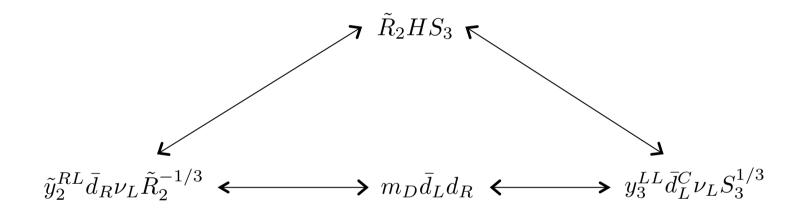


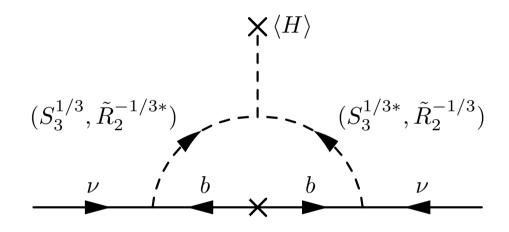
 $y_3^{LL} \bar{d}_L^C
u_L S_3^{1/3}$





$$H \equiv (\mathbf{1}, \mathbf{2}, -1/2)$$





IMPORTANT ISSUES

WHAT HAPPENED WITH THE LQ DIQUARK COUPLINGS?

LQ MASSES ARE FREE PARAMETERS...

 $\tilde{y}_2^{RL}, y_1^{LL}, y_3^{LL}$ ARE ALL *A PRIORI* UNKNOWN MATRICES...

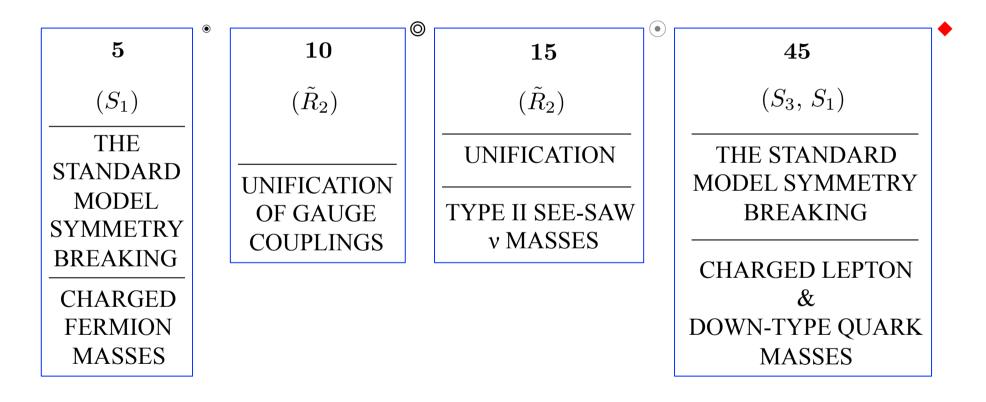
RECENT DEVELOPMENTS[©]

"A testable radiative neutrino mass model without additional symmetries and explanation for the $b \rightarrow s\ell^+\ell^-$ anomaly"

[©]K. Cheung, T. Nomura, H. Okada, arXiv:1610.02322.

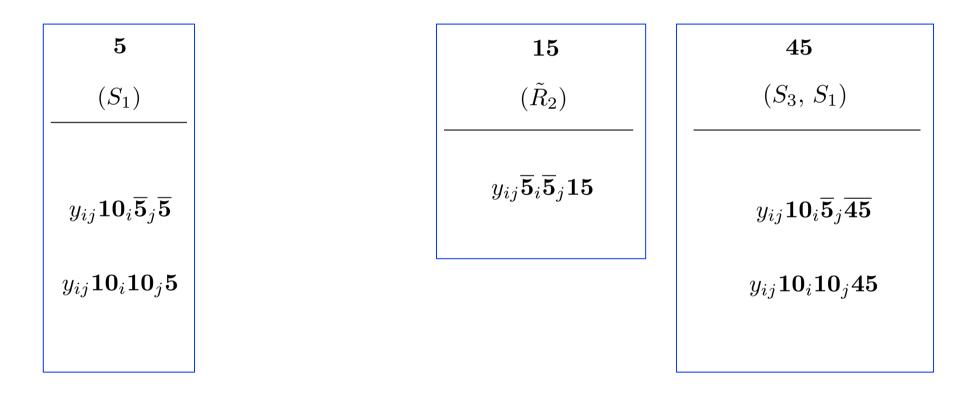
SCALAR LQs	SU(5)
$S_3 \equiv (\overline{3}, 3, 1/3)$	$\overline{45}$
$R_2 \equiv (3, 2, 7/6)$	$\overline{45}$
$\tilde{R}_2 \equiv (3, 2, 1/6)$	10, 15
$\tilde{S}_1 \equiv (\overline{3}, 1, 4/3)$	45
$S_1 \equiv (\overline{3}, 1, 1/3)$	$\overline{5}, \overline{45}$
$\bar{S}_1 \equiv (\bar{3}, 1, -2/3)$	10

SCALAR REPRESENTATIONS IN *SU*(5):



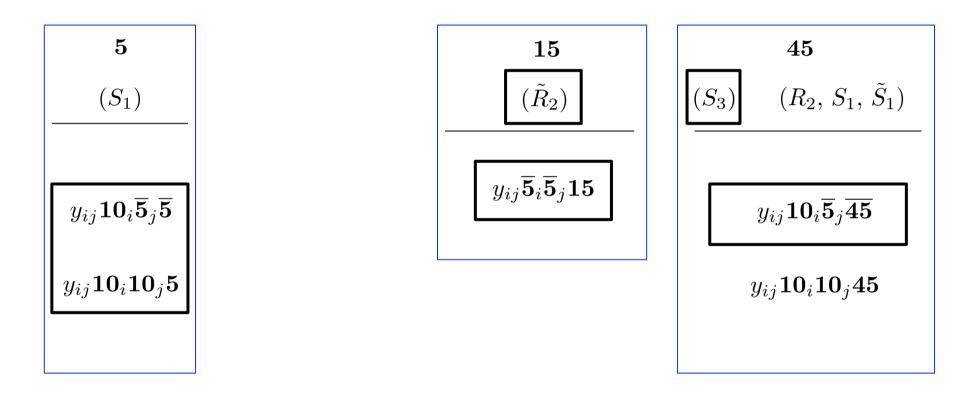
[●]H. Georgi, S. L. Glashow, Phys. Rev. Lett. 32 (1974). [◎]H. Murayama, T. Yanagida, Mod. Phys. Lett. A7 (1992). [●]I. Doršner, P. Fileviez Perez, Nucl. Phys. B723 (2005). [◆]H. Georgi, C. Jarlskog., Phys. Lett. B86 (1979).

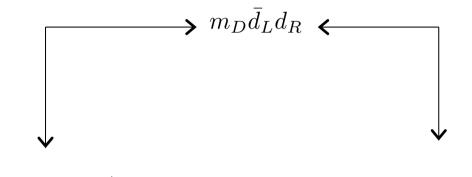
SCALAR REPRESENTATIONS IN *SU*(5):



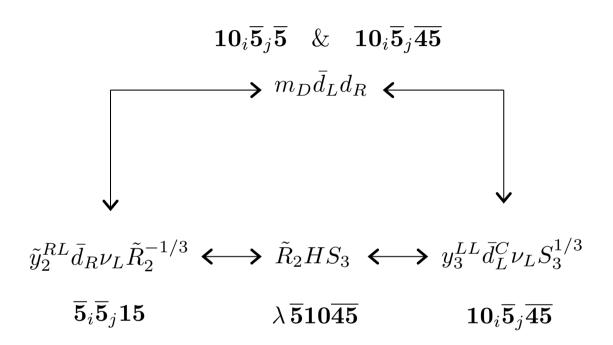
 $\mathbf{10}_i \& \overline{\mathbf{5}}_i \ (i = 1, 2, 3)$ COMPRISE THE STANDARD MODEL FERMIONS

A POSSIBLE *SU*(5) SET-UP:





 $\tilde{y}_2^{RL} \bar{d}_R \nu_L \tilde{R}_2^{-1/3} \longleftrightarrow \tilde{R}_2 H S_3 \longleftrightarrow y_3^{LL} \bar{d}_L^C \nu_L S_3^{1/3}$



 $\lambda \equiv \text{dimensionful parameter}$

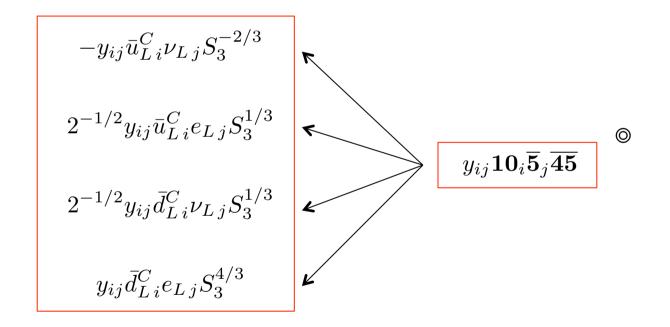


$y_3^{LL} \bar{d}_L^C \nu_L S_3^{1/3}$

 $y_{ij} \mathbf{10}_i \overline{\mathbf{5}}_j \overline{\mathbf{45}}$

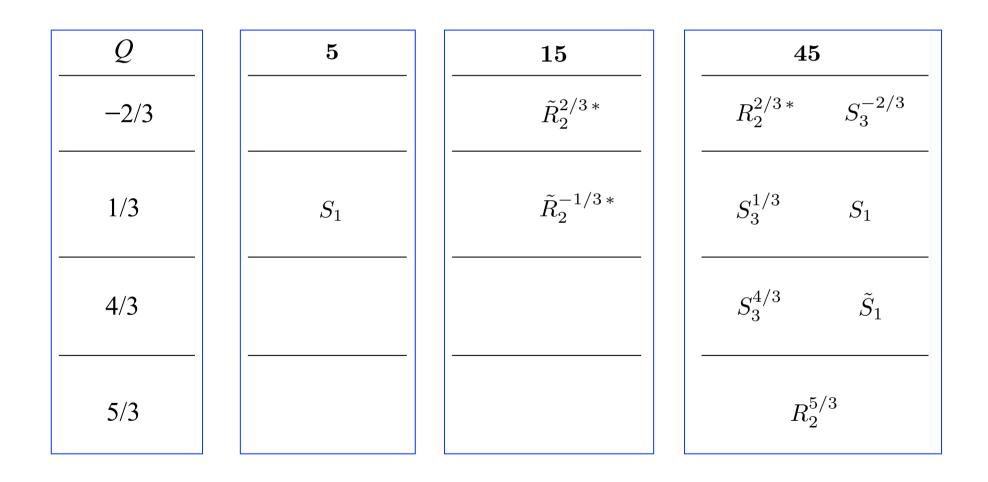
<u>p DECAY</u>

S_3 LEPTOQUARK MULTIPLET COULD BE LIGHT IF NEEDED \ldots

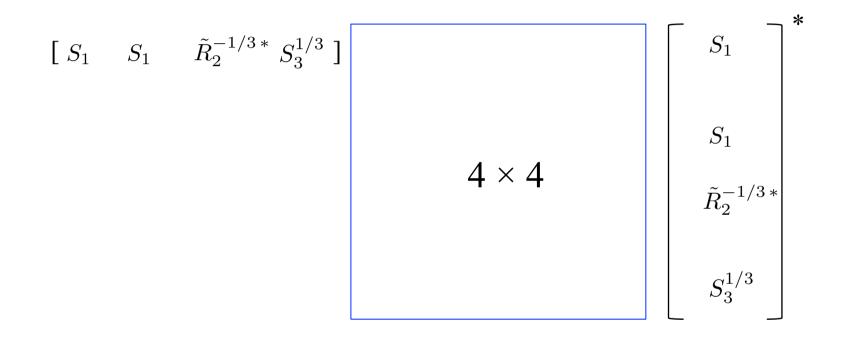


[©]I. Doršner, S. Fajfer, N. Košnik, Phys. Rev. D 86, 015013 (2012).

<u>p DECAY</u>



p **DECAY**



<u>p DECAY</u>

$$\begin{bmatrix} S_1 & S_1 & \tilde{R}_2^{-1/3*} & S_3^{1/3} \end{bmatrix} \begin{bmatrix} & & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & \\ & & & & & & \\ & & & & & & \\ & & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & \\ & & & & & \\ & & & & \\ & & & & & \\ & & &$$

$$\begin{bmatrix} \tilde{R}_{2}^{-1/3} * S_{3}^{1/3} \end{bmatrix} \begin{bmatrix} m_{11}^{2} & m_{12}^{2} \\ m_{12}^{2} & m_{22}^{2} \end{bmatrix} \begin{bmatrix} \tilde{R}_{2}^{-1/3} * \\ S_{3}^{1/3} \end{bmatrix}^{*}$$

$$(m_{\nu})_{ij} = \frac{3s_{\theta}c_{\theta}}{16\pi^2} \sum_{k=d,s,b} m_k [B_0(0, m_k^2, m_1^2) - B_0(0, m_k^2, m_2^2)] \{y_{ik}y'_{jk} + y_{jk}y'_{ik}\}$$

$$B_0(0, m_k^2, m_1^2) - B_0(0, m_k^2, m_2^2) = \frac{m_2^2 [\ln m_2^2 - \ln m_k^2]}{m_2^2 - m_k^2} - \frac{m_1^2 [\ln m_1^2 - \ln m_k^2]}{m_1^2 - m_k^2}$$

B_0 – Passarino-Veltman function

$$(m_{\nu})_{ij} = \frac{3s_{\theta}c_{\theta}}{16\pi^2} \sum_{k=d,s,b} m_k [B_0(0, m_k^2, m_1^2) - B_0(0, m_k^2, m_2^2)] \{y_{ik}y'_{jk} + y_{jk}y'_{ik}\}$$

$$y_{ij}\overline{\mathbf{5}}_i\overline{\mathbf{5}}_j\mathbf{15}$$
 $y'_{ij}\mathbf{10}_i\overline{\mathbf{5}}_j\overline{\mathbf{45}}$

$$(m_{\nu})_{ij} = \frac{3s_{\theta}c_{\theta}}{16\pi^2} \sum_{k=d,s,b} m_k [B_0(0, m_k^2, m_1^2) - B_0(0, m_k^2, m_2^2)] \{y_{ik}y'_{jk} + y_{jk}y'_{ik}\}$$

$$y_{ij}\overline{\mathbf{5}}_i\overline{\mathbf{5}}_j\mathbf{15}$$
 $y' \sim (M_e^T - M_d)/v_{45}$

 M_e – charged lepton mass matrix

 M_d – down-type quark mass matrix

A LIST OF BENEFITS

p DECAY CONSTRAINTS CAN BE ACCOMMODATED

© RELEVANT YUKAWA COUPLING MATRICES ARE RELATED TO FERMION MASSES AND/OR POSSESS ADDITIONAL SYMMETRY. THIS NOT ONLY REDUCES THE TOTAL NUMBER OF PARAMETERS BUT HELPS RELATE LEPTOQUARK DECAY PATTERNS TO NEUTRINO MASSES...

LQ MASSES COULD BE CONSTRAINED THROUGH THE GAUGE COUPLING UNIFICATION...

[©]P. Fileviez Perez, T. Han, Gui-Yu Huang, T. Li, K. Wang, Phys. Rev. D 78, 071301, (2008).

CONCLUSIONS

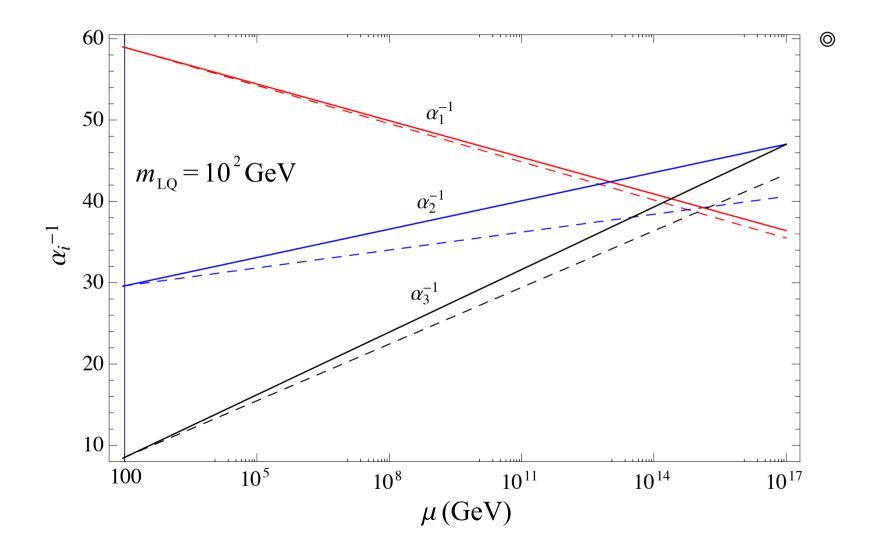
SU(5) CAN ACCOMMODATE WITH EASE THE ONE-LOOP NEUTRINO MASS MECHANISM THAT IS BASED ON THE LEPTOQUARK MULTIPLET MIXING.

THE USE OF *SU*(5) CAN INCREASE PREDICTIVITY OF THE SET-UP. THIS COULD ESPECIALLY BE REFLECTED IN THE DECAY PATTERNS OF THE RELEVANT LEPTOQUARK MULTIPLETS.

THANK YOU

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STANDARD MODEL + (2×10)



[©]H. Murayama, T. Yanagida, Mod. Phys. Lett. A7 (1992).