

# Light window for right-handed neutrinos in the Left-Right model

Goran Popara  
(with F. Nesti and M. Nemevšek)

Ruder Bošković Institute

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Existence of neutrino masses implied by the discovery of neutrino oscillations (Super-Kamiokande & SNO).

Important questions:

- ▶ Dirac/Majorana nature of neutrinos,
- ▶ neutrino mass generation mechanism.

# Introduction

Attractive explanation of neutrino masses (and their smallness) is the *seesaw mechanism*, where neutrinos are Majorana fermions.

Possible extension of SM is the *Left-Right symmetric model* (LRSM):

- ▶ restores parity,
- ▶ naturally embeds the seesaw mechanism.

Important consequence of Majorana neutrinos is *lepton number violation* (LNV),  $\Delta L = 2$ .

# Left-right symmetric model

J. C. Pati, A. Salam, PRD **10** (1974); **11** (1975); R. N. Mohapatra, PRD **11** (1975)  
G. Senjanović, R. N. Mohapatra, PRD **12** (1975); G. Senjanović, PRL **44** (1980) ...

Gauge group:

$$\mathcal{G}_{LR} = SU(2)_L \times SU(2)_R \times U(1)_{B-L}$$
$$\Rightarrow W_{L,R} \quad Z_{L,R} \quad \gamma$$

Matter fields:

$$Q_{L,i} = \begin{pmatrix} u_L \\ d_L \end{pmatrix}_i \sim \left( \mathbf{2}, \mathbf{1}, \frac{1}{3} \right) \quad Q_{R,i} = \begin{pmatrix} u_R \\ d_R \end{pmatrix}_i \sim \left( \mathbf{1}, \mathbf{2}, \frac{1}{3} \right)$$
$$\psi_{L,i} = \begin{pmatrix} \nu_L \\ l_L \end{pmatrix}_i \sim (\mathbf{2}, \mathbf{1}, -1) \quad \psi_{R,i} = \begin{pmatrix} N_R \\ l_R \end{pmatrix}_i \sim (\mathbf{1}, \mathbf{2}, -1)$$

# Left-right symmetric model

Scalar sector:

$$\Phi = \begin{pmatrix} \phi_1^0 & \phi_2^+ \\ \phi_1^- & \phi_2^0 \end{pmatrix} \sim (\mathbf{2}, \mathbf{2}, 0)$$

$$\Delta_{L,R} = \begin{pmatrix} \Delta^+/\sqrt{2} & \Delta^{++} \\ \Delta^0 & -\Delta^+/\sqrt{2} \end{pmatrix}_{L,R} \sim (\mathbf{3}, \mathbf{1}, 2), (\mathbf{1}, \mathbf{3}, 2)$$

Symmetry breaking pattern:

$$\mathcal{G}_{LR} \xrightarrow[\langle \Delta_L \rangle = 0]{\langle \Delta_R \rangle \neq 0} SU(2)_L \times U(1) \xrightarrow{\langle \Phi \rangle \neq 0} U(1)_{\text{em}}$$

$$Q_{\text{em}} = I_{3L} + I_{3R} + \frac{B - L}{2}$$

# Left-right symmetric model

Crucial ingredient — *Majorana nature of neutrinos*.

Smallness of neutrino mass is connected to the scale of new physics ( $M_{W_R}$ ):

$$m_{\nu_l} \sim \frac{m_l^2}{M_{W_R}}.$$

# Constraints

Constraints from low-energy experiments:

- ▶  $K^0 - \bar{K}^0$  and  $B_{d,s}^0 - \bar{B}_{d,s}^0$  oscillations

Y. Zhang, H. An, X. Ji, R. N. Mohapatra, Nucl. Phys. B **802** (2008);

S. Bertolini, A. Maiezza, F. Nesti, Phys. Rev. D **89** (2014)

- ▶ CP-violating processes ( $\varepsilon$ ,  $\varepsilon'$ )

S. Bertolini, J. O. Eeg, A. Maiezza, F. Nesti, Phys. Rev. D **86** (2012);

S. Bertolini, A. Maiezza, F. Nesti, Phys. Rev. D **88** (2013)

- ▶  $n$ EDM

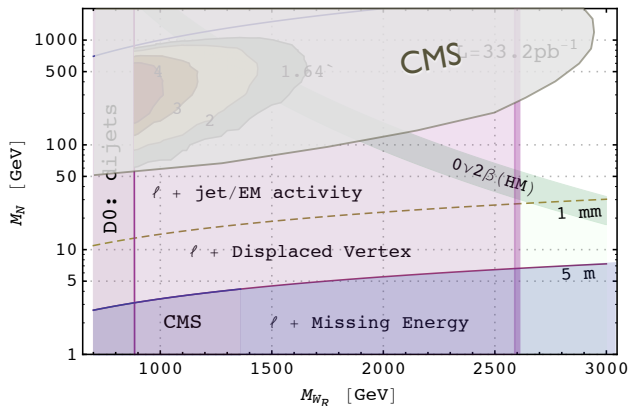
A. Maiezza, M. Nemevšek, Phys. Rev. D **90** (2014)

$$\Rightarrow M_{W_R} \gtrsim 3 \text{ TeV}$$

**Neutrino nautre is observable at LHC!**

# Constraints: Light window

M. Nemevšek, F. Nesti, G. Senjanović, Y. Zhang, Phys. Rev. D **83** (2011)



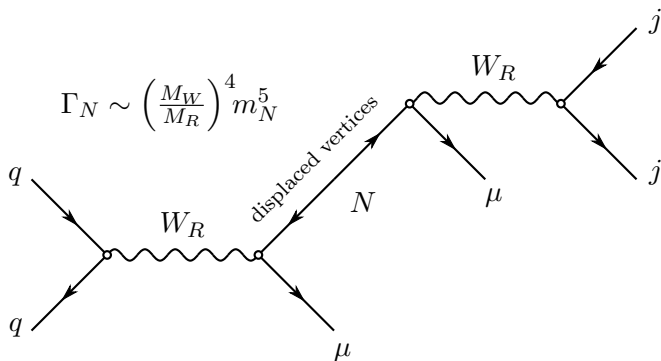
$$4 \text{ TeV} \lesssim M_{W_R} \lesssim 7 \text{ TeV} \quad \text{and} \quad 10 \text{ GeV} \lesssim m_N \lesssim 400 \text{ GeV}$$



# Keung-Senjanović process

W.-Y. Keung, G. Senjanović, Phys. Rev. Lett. **50** (1983)

$$pp \rightarrow \mu^\pm \mu^\pm jj \quad (\text{LNV!})$$



# Keung-Senjanović process

Important features of Keung-Senjanović (KS) process:

- ▶ lepton number violation (not present in SM),
- ▶ *displaced vertices* (long-lived  $N$ ) — helpful in eliminating the background,
- ▶ high-energy analogue to  $0\nu 2\beta$ .

Although clean in principle, presents some real-world challenges:

- ▶ background
- ▶ detector effects (e. g. highly boosted products of  $N$  decay become collimated).

# Monte Carlo simulation

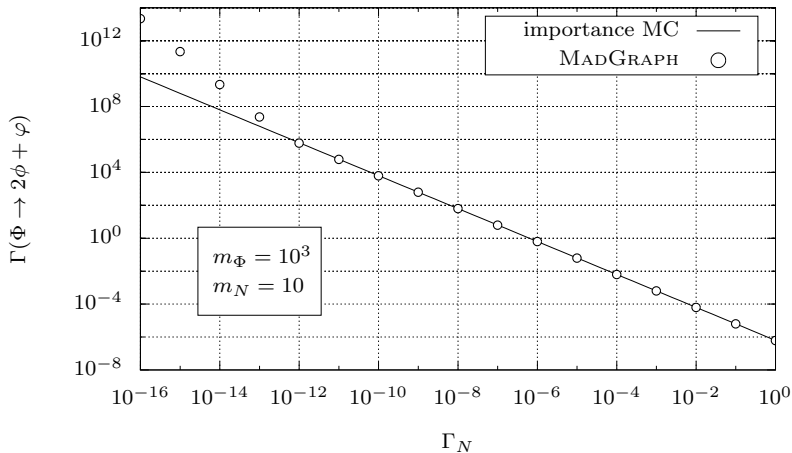
Simulation of signal and background involves several steps:

1. model definition (FeynRules),
2. event generation (MadGraph),
3. hadronization (Pythia),
4. detector simulation (Delphes),
5. analysis, cuts (MadAnalysis).

Narrow  $N$  resonance causes numerical instabilities in the event generation step!

# Monte Carlo simulation: Narrow width problem

Simple test with scalars:  $\Phi \rightarrow \phi N, N \rightarrow \phi\phi$ .



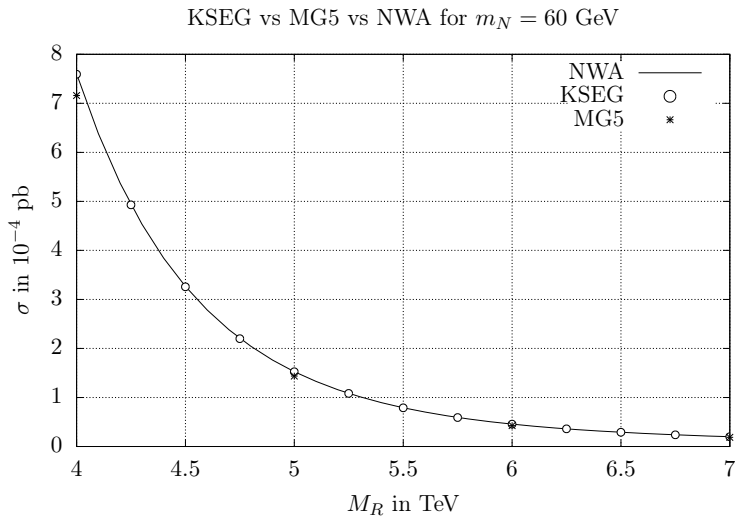
# Monte Carlo simulation: Custom event generator

Solution – custom event generator (KSEG):

- ▶ recursive phase space decomposition (integration over resonances),
- ▶ importance sampling from Breit-Wigner distribution to eliminate narrow  $N$  peak,
- ▶ unweighted events output to LHE file (for further processing).

# Monte Carlo simulation: Custom event generator

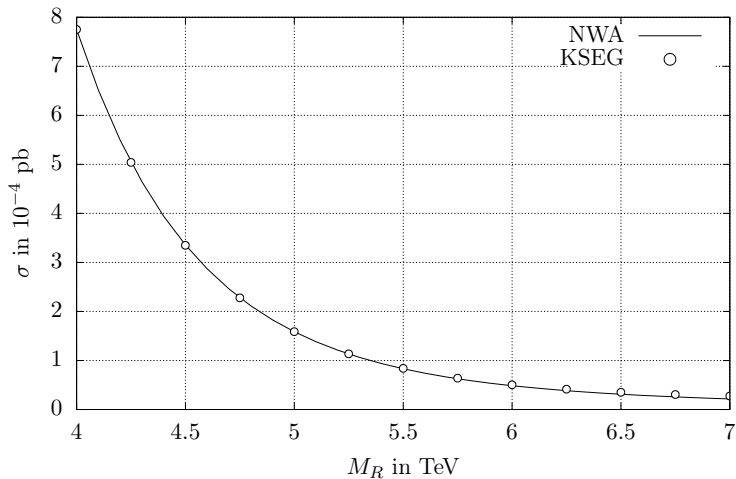
Cross-check using narrow width approximation:



# Monte Carlo simulation: Custom event generator

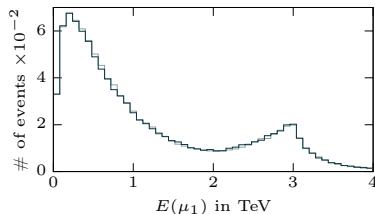
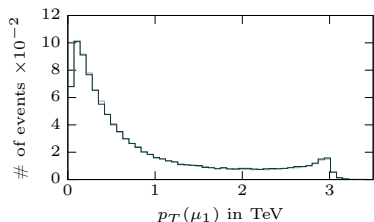
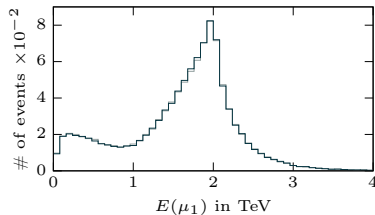
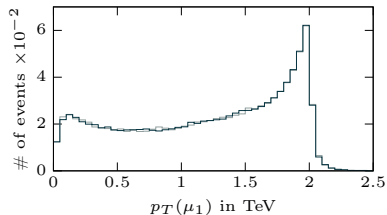
Cross-check using narrow width approximation:

KSEG vs NWA for  $m_N = 10$  GeV



# Monte Carlo simulation: Custom event generator

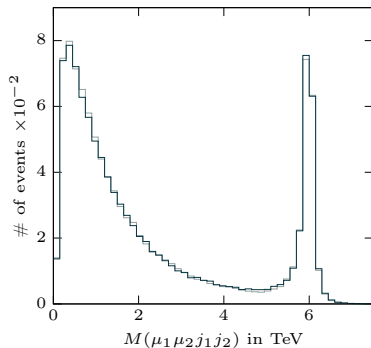
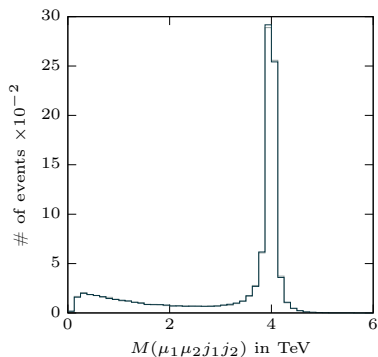
Transverse momentum and energy distributions (KSEG & MG5) of the prompt muon for  $m_N = 80$  GeV and  $M_R = 4$  TeV (upper panel) and  $M_R = 6$  TeV (lower panel):





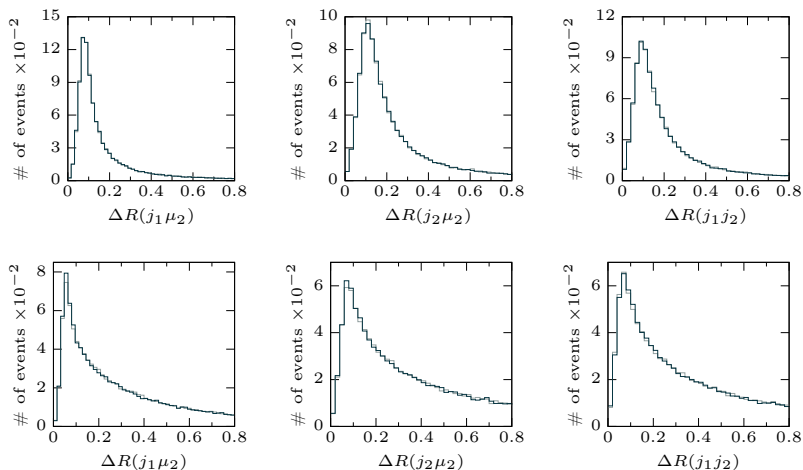
# Monte Carlo simulation: Custom event generator

Invariant mass of the muons and jets for  $m_N = 80$  GeV and  $M_R = 4$  TeV (left) and  $M_R = 6$  TeV (right):



# Monte Carlo simulation: Custom event generator

$\Delta R$  of various pairs for  $m_N = 80$  GeV and  $M_R = 4$  TeV (upper panel) and  $M_R = 6$  TeV (lower panel):



# Outlook

MadGraph instability problem in Keung-Senjanović process is solved (for the relevant portion of parameter space).

Hadronization and detector simulation ongoing.

Things to be done:

1. consider background,
2. place/optimize cuts (displaced vertices interesting here),
3. sensitivity assessment, limits on the model,
4. prospects for future colliders.

Thank you!

Two kinds of LR symmetries, imposing restrictions on Yukawa matrices:

$$\mathcal{P} : \begin{cases} Q_L \leftrightarrow Q_R \\ \Phi \rightarrow \Phi^\dagger \end{cases} \Rightarrow Y = Y^\dagger, \quad \mathcal{C} : \begin{cases} Q_L \leftrightarrow (Q_R)^c \\ \Phi \rightarrow \Phi^T \end{cases} \Rightarrow Y = Y^T.$$

$\mathcal{C}$  has an advantage — it can be gauged (involves spinors with same final chirality).

A. Maiezza, M. Nemevšek, F. Nesti, G. Senjanović, Phys. Rev. D **82** (2010)