

## Nonperturbative thermodynamics of multi-Higgs models

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In collaboration with:

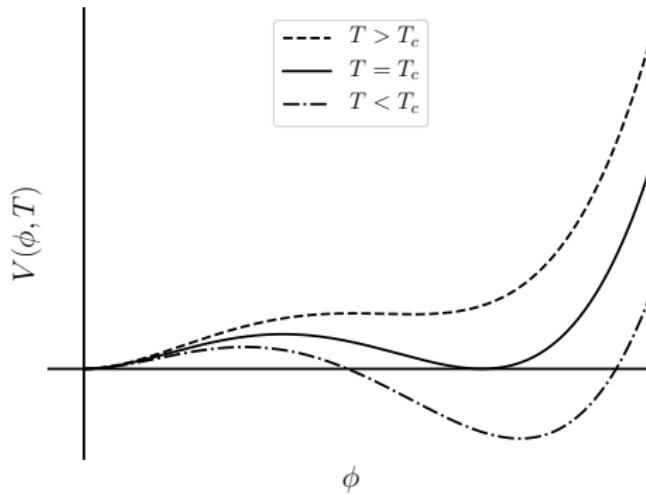
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Based on 1904.01329, 1903.11604

# First-order phase transitions in QFT

- $T = 0$  effective potential + thermal corrections:



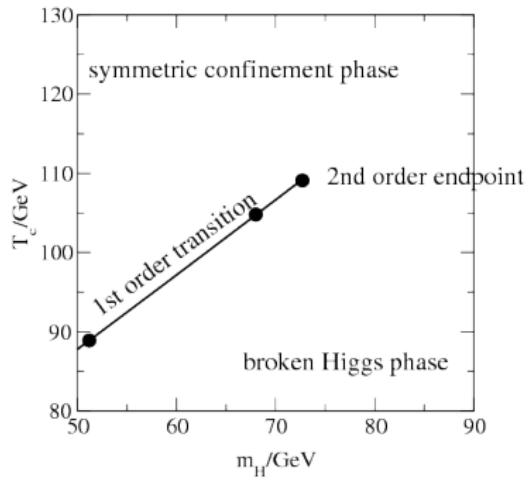
- Energy minima separated by a potential barrier  $\rightarrow$  latent heat from tunneling

# First order electroweak phase transition?

NO phase transition in the SM (smooth crossover)!

Need significant modifications  
to EW scale dynamics from BSM physics:

- ▶ Through radiative corrections to the Higgs (often  $\mathcal{O}(1)$  couplings)?
- ▶ Multiple phase transitions before settling to the EW minimum?
- ▶ Non-renormalizable operators?



Source: hep-ph/0010275

# Why not perturbation theory?

- Infrared problem: light bosons are nonperturbative

$$\text{expansion parameter} \sim g^2 n_b(m) = \frac{g^2}{e^{m/T} - 1} \xrightarrow{m \ll T} \frac{g^2 T}{m}$$

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Examples:

- ▶ Perturbation theory underestimates transition strength by  $\sim 50\%$  in MSSM      Laine et al. '12
- ▶ Perturbative effective potential always has a barrier

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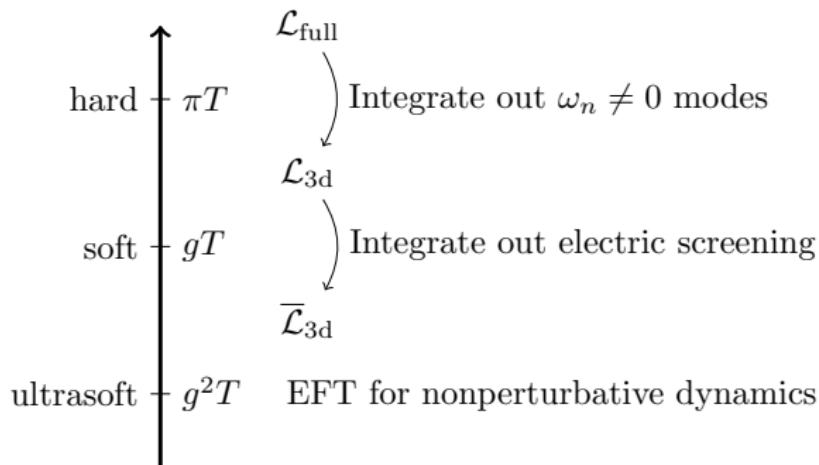
- What are the heavy modes?  
→ Thermal fluctuations of inverse length  $\lesssim \pi T$ :

$$\phi(\tau, \mathbf{x}) = T \sum_n \tilde{\phi}(\mathbf{p}) e^{i\omega_n \tau}, \quad \omega_n = \begin{cases} 2\pi n T & \text{bosons} \\ (2n + 1)\pi T & \text{fermions} \end{cases}$$

Matsubara decomposition

# High- $T$ effective theories

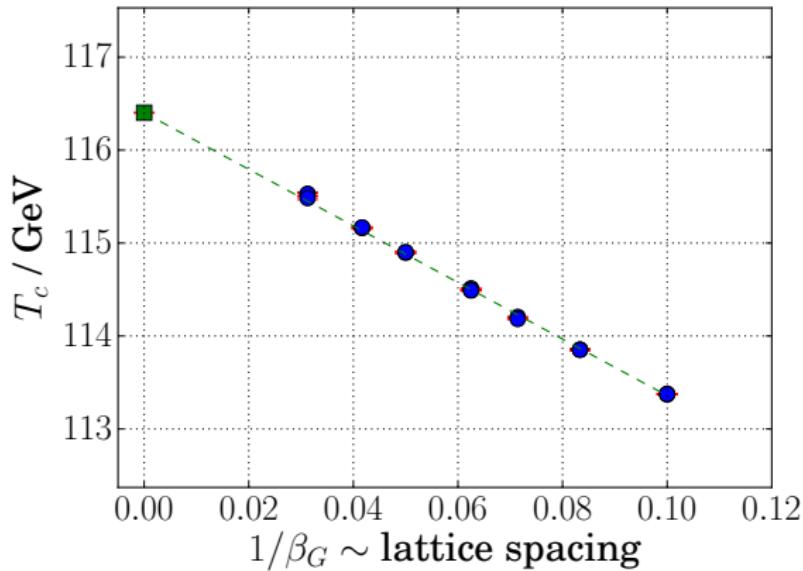
- Temperature decouples  $\rightarrow$  IR theory is three dimensional and easy to simulate
- Easily applied to (perturbative) BSM theories



# Lattice simulations provide accurate results

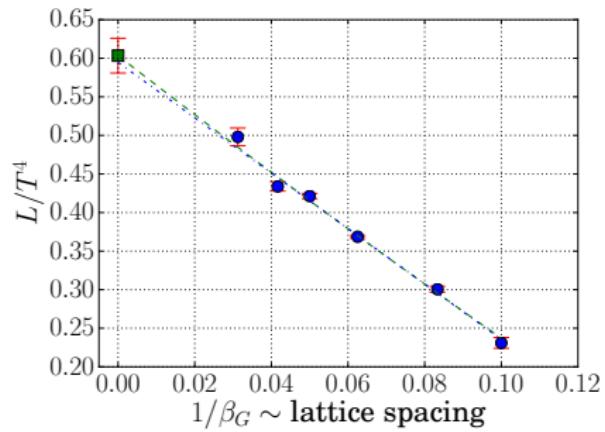
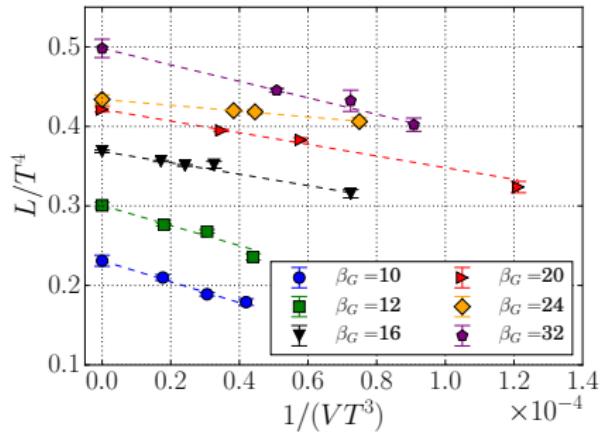
Example: phase transition with two Higgs doublets

Kainulainen et al. 1904.01329

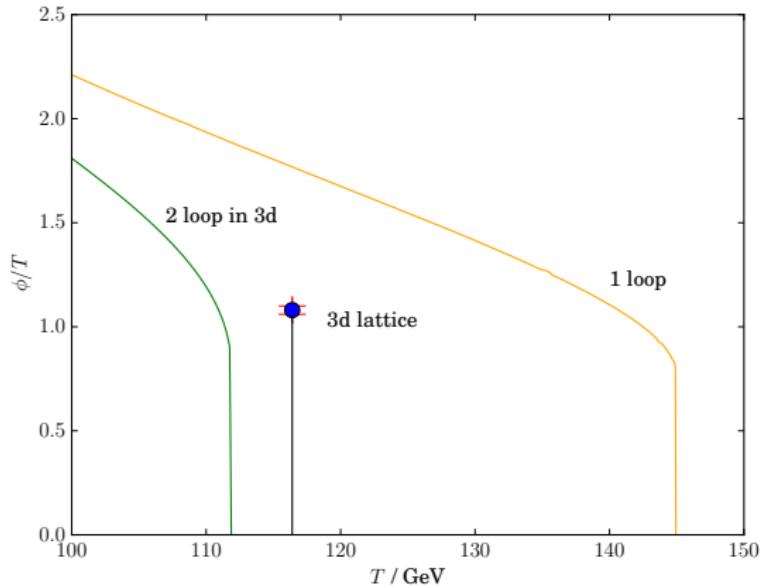


# Continuum extrapolation for latent heat

Measurable from discontinuity of condensates  $\langle \phi^\dagger \phi \rangle$ ,  $\langle (\phi^\dagger \phi)^2 \rangle$ , ...

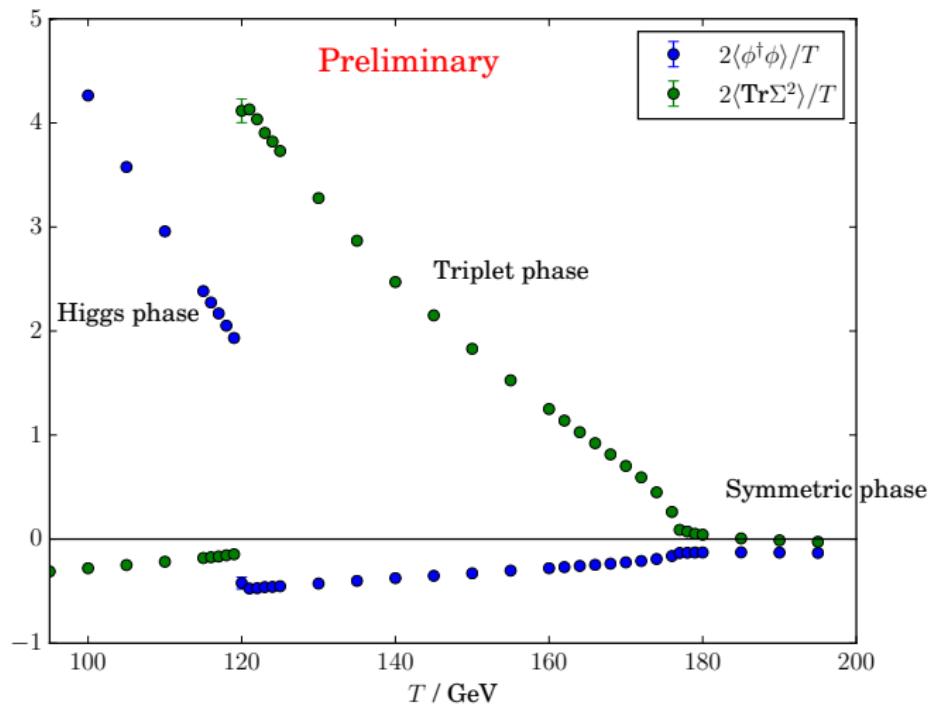


# Comparison with perturbation theory



# Multiple phase transitions?

SM + SU(2) triplet scalar (finite lattice):



# Model independent study?

EFT for SM + any sufficiently heavy scalar:

$$\begin{aligned}\mathcal{L}_{3d} = & \frac{1}{4} F_{ij}^a F_{ij}^a + (D_i \phi)^\dagger (D_i \phi) + m_3^2 \phi^\dagger \phi + \lambda_3 (\phi^\dagger \phi)^2 \\ & + \text{neglected higher-order operators}\end{aligned}$$

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+ neglected higher-order operators

Nonperturbative phase diagram known Kajantie et al. '96

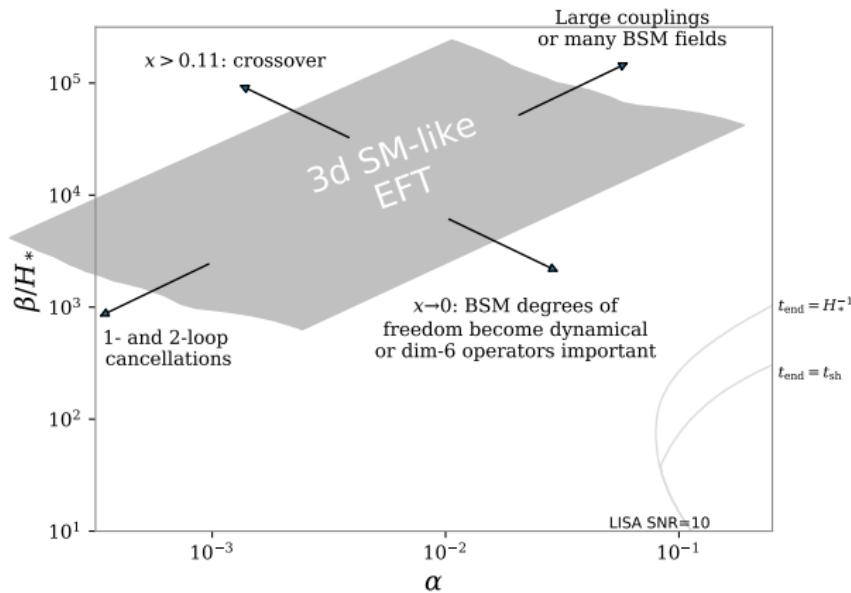
$$x \equiv \frac{\lambda_3}{g_3^2}, \quad y \equiv \frac{m_3^2}{g_3^4} \approx 0 \text{ at } T_c$$

- ▶ First order transition for  $0 < x < 0.11$
- ▶ Crossover for  $x > 0.11$

# Gravitational waves from the SM-like EFT

Nonperturbative nucleation rate known

Moore & Rummukainen '01



No observable gravitational wave signals from any theory that looks like SM in the infrared at high- $T$ !

arXiv:1903.11604

# Conclusions

- ▶ New physics contributions can turn the EW phase transition into first order
- ▶ Calculations at high temperature can be simplified with effective theory methods
- ▶ Lattice methods can provide quantitative results, and can be used to estimate the performance of perturbation theory

## Backup: Properties of the effective theory

Example: Two Higgs doublet model.

$$V(\phi_1, \phi_2) = m_1^2 \phi_1^\dagger \phi_1 + m_2^2 \phi_2^\dagger \phi_2 + m_{12}^2 \phi_1^\dagger \phi_2 + \dots + \lambda_3 (\phi_1^\dagger \phi_1)(\phi_2^\dagger \phi_2) + \dots$$

The resulting EFT:

$$\begin{aligned}\bar{\mathcal{L}}_{3d} = & \frac{1}{4} (F_{rs})^2 + (D_r \phi_1)^\dagger (D_r \phi_1) + (D_r \phi_2)^\dagger (D_r \phi_2) + \overline{m}_1^2(T) \phi_1^\dagger \phi_1 \\ & + \overline{m}_2^2(T) \phi_2^\dagger \phi_2 + \overline{m}_{12}^2(T) \phi_1^\dagger \phi_2 + \dots + \overline{\lambda}_3(T) (\phi_1^\dagger \phi_1)(\phi_2^\dagger \phi_2) + \dots\end{aligned}$$

- ▶ Super-renormalizable → **exact** relations to lattice parameters.
- ▶ Very accurate in the weak coupling regime. Error in the SM is  $\sim 1\%$ .

## Backup: Finding the transition point from histograms

- Find precise  $T_c$  by reweighting the histograms to equal weight.

